



Fig. 6. Same as Fig. 5, except that emphasis is on next-nearest-neighbor pairs. (a) Undeformed state; all three  $\langle 100 \rangle$  NNN directions contain AA and BB pairs only. (b) to (d) Alternative arrangements after  $(\bar{1}12) [1\bar{1}1]$  slip, showing AB pairs (double bars) induced in all three  $\langle 100 \rangle$  directions.

Since  $\Delta N_{BB}$  occurs in  $[111]$ ,  $[\bar{1}\bar{1}1]$  and  $[11\bar{1}]$ , we have from eqn. (1)

$$\begin{aligned} \cos^2 \varphi &= (\cos^2 \varphi_{[111]} + \cos^2 \varphi_{[\bar{1}\bar{1}1]} + \cos^2 \varphi_{[11\bar{1}]}) \\ &= -\frac{2}{3}(\alpha_1 \alpha_2 - \alpha_2 \alpha_3 + \alpha_3 \alpha_1), \end{aligned} \quad (15)$$

which means that the  $[1\bar{1}1]$  slip direction is the effective direction of unlike pairs, which is a hard direction of magnetization for FeCo. In the general case, where the slip direction of slip system  $i$  has the direction cosines  $d_{1i}, d_{2i}, d_{3i}$ ,

$$\begin{aligned} \cos^2 \varphi_i &= -2(d_{1i}d_{2i}\alpha_1\alpha_2 + d_{2i}d_{3i}\alpha_2\alpha_3 \\ &\quad + d_{3i}d_{1i}\alpha_3\alpha_1) \\ &\equiv -2f_i(\alpha_1, \alpha_2, \alpha_3). \end{aligned} \quad (16)$$

The insertion of eqns. (14) and (16) into eqn. (1) then leads to

$$E = -\frac{1}{2}Nlp_0p's^2 \sum_i |S_i| f_i(\alpha_1, \alpha_2, \alpha_3). \quad (17)$$

(b) Short-range order case

For the short-range order case, the two nearest-neighbor bonds associated with each  $\langle 111 \rangle$  direction (in a unit cell) contain a total of  $(1-\sigma)/2$  BB pairs in the undeformed state (see Appendix). After slip,  $\sigma=0$  for the 3 $\langle 111 \rangle$  directions other than the slip direction. Hence the gain in BB pairs in each of these directions is  $\sigma/2$ , or  $\sigma/2a^2\sqrt{6}$  per

TABLE III: VALUES OF  $\varepsilon$  AND  $d$  (REFERRED TO CUBIC AXES) FOR THE TWELVE  $\{112\} \langle 111 \rangle$  SLIP SYSTEMS

No. of slip system	Slip plane	Slip direction	$2\varepsilon_{xx}$	$2\varepsilon_{yy}$	$2\varepsilon_{zz}$	$4\varepsilon_{yz}$	$4\varepsilon_{zx}$	$4\varepsilon_{xy}$	$3d_1d_2$	$3d_2d_3$	$3d_3d_1$
1	(11 $\bar{2}$ )	111	$S_1$	$S_1$	$-2S_1$	$-S_1$	$-S_1$	$2S_1$	1	1	1
2	(1 $\bar{2}$ 1)	111	$S_2$	$-2S_2$	$S_2$	$-S_2$	$2S_2$	$-S_2$	1	1	1
3	( $\bar{2}$ 11)	111	$-2S_3$	$S_3$	$S_3$	$2S_3$	$-S_3$	$-S_3$	1	1	1
4	(112)	11 $\bar{1}$	$S_4$	$S_4$	$-2S_4$	$S_4$	$S_4$	$2S_4$	1	-1	-1
5	( $\bar{1}$ 21)	11 $\bar{1}$	$-S_5$	$2S_5$	$-S_5$	$-S_5$	$2S_5$	$S_5$	1	-1	-1
6	(2 $\bar{1}$ 1)	11 $\bar{1}$	$2S_6$	$-S_6$	$-S_6$	$2S_6$	$-S_6$	$S_6$	1	-1	-1
7	(1 $\bar{1}$ 2)	$\bar{1}$ 11	$-S_7$	$-S_7$	$2S_7$	$S_7$	$-S_7$	$2S_7$	-1	1	-1
8	(211)	$\bar{1}$ 11	$-2S_8$	$S_8$	$S_8$	$2S_8$	$S_8$	$S_8$	-1	1	-1
9	(12 $\bar{1}$ )	$\bar{1}$ 11	$-S_9$	$2S_9$	$-S_9$	$S_9$	$2S_9$	$-S_9$	-1	1	-1
10	(21 $\bar{1}$ )	$\bar{1}$ 11	$2S_{10}$	$-S_{10}$	$-S_{10}$	$2S_{10}$	$S_{10}$	$-S_{10}$	-1	-1	1
11	(121)	1 $\bar{1}$ 1	$S_{11}$	$-2S_{11}$	$S_{11}$	$S_{11}$	$2S_{11}$	$S_{11}$	-1	-1	1
12	( $\bar{1}$ 12)	1 $\bar{1}$ 1	$-S_{12}$	$-S_{12}$	$2S_{12}$	$-S_{12}$	$S_{12}$	$2S_{12}$	-1	-1	1

TABLE IV: SUMMARY OF RESULTS BASED ON  $\{112\} \langle 111 \rangle$  SLIP

Rolling plane	Rolling direction	Active slip systems	$ S_i $	$E_{NN}$	Easy* axis
(001)	[ $\bar{1}$ 10]	7,12	$r/2$	$\left(\frac{E_1 r}{6}\right) \alpha_1 \alpha_2$	TD
(115)	[ $\bar{1}$ 10]	4,7,8,11,12	$ S_4  = 2r/27$ $ S_8  =  S_{11}  = 6r/27$ $ S_7  =  S_{12}  = 21r/54$	$\left(\frac{E_1 r}{162}\right) (31\alpha_1 \alpha_2 + 2\alpha_2 \alpha_3 + 2\alpha_3 \alpha_1)$	TD
(112)	[ $\bar{1}$ 10]	4,7,8,11,12	$r/3$	$\left(\frac{E_1 r}{18}\right) (3\alpha_1 \alpha_2 + \alpha_2 \alpha_3 + \alpha_3 \alpha_1)$	RPN
(111)	[ $\bar{1}$ 10]	4,7,8,11,12	$ S_4  = 2r/3$ $ S_8  =  S_{11}  = 2r/9$ $ S_7  =  S_{12}  = 7r/18$	$\left(\frac{E_1 r}{54}\right) (5\alpha_1 \alpha_2 + 6\alpha_2 \alpha_3 + 6\alpha_3 \alpha_1)$	RPN
(110)	[ $\bar{1}$ 10]	1,4,7,12	$r/2$	0	—

\* Relative among the three symmetry directions of rolled strip.  
TD-transverse direction, RPN-rolling plane normal.

unit  $\{112\}$  area. To this quantity we multiply by a factor  $(p'|S|/d = p'|S|\sqrt{6}/a)$  as before, leading to the expression

$$\Delta N_{BB} = p'|S| \left( \frac{\sqrt{6}}{a} \right) \left( \frac{\sigma}{2a^2\sqrt{6}} \right) = \frac{1}{4} N p' \sigma |S|. \quad (18)$$

Insertion of eqns. (16) and (18) into eqn. (1) then leads to

$$E = -\frac{1}{2} N l p' \sigma \sum_i |S_i| f_i(\alpha_1, \alpha_2, \alpha_3) \quad (19)$$

for the SRO case.

Finally, by combining the LRO and SRO expressions of eqns. (17) and (19), we obtain

$$E = -\frac{1}{2} E_1 \sum_i |S_i| f_i(\alpha_1, \alpha_2, \alpha_3), \quad (20)$$

where  $E_1 = N l p' (p_0 s^2 + \sigma)$  as before. Equation (20) is the expression for the slip-induced anisotropy energy developed by  $\{112\} \langle 111 \rangle$  slip and based on NN interactions.

### (c) Applications to rolling

As in the treatment of  $\{110\} \langle 111 \rangle$  slip in Part 2 above, eqn. (20) has been applied to rolling of (001)[ $\bar{1}$ 10], (115)[ $\bar{1}$ 10], (112)[ $\bar{1}$ 10], (111)[ $\bar{1}$ 10] and (110)[ $\bar{1}$ 10] orientations. For the convenience of calculations, the appropriate parameters (similar to